

Theoretical Description of Coherent Doublon Creation via Lattice Modulation Spectroscopy

Andreas Dirks* and Karlis Mikelsons

Department of Physics, Georgetown University, Washington, DC 20057, USA

H. R. Krishnamurthy

*Centre for Condensed Matter Theory, Department of Physics,
Indian Institute of Science, Bangalore 560012, India, and*

Jawaharlal Nehru Centre for Advanced Scientific Research, Bangalore 560064, India

J. K. Freericks

*Department of Physics, Georgetown University, Washington, DC 20057, USA and
Kavli Institute for Theoretical Physics, Santa Barbara, CA 93106, USA*

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Using a recently developed strong-coupling method, we present a comprehensive theory for doublon production processes in modulation spectroscopy of a system of ultracold fermionic atoms in an optical lattice. The theoretical predictions compare very well to the experimental time traces of doublon production. For experimentally feasible conditions, we provide a quantitative prediction for the presence of a nonlinear "two-photon" excitation at strong modulation amplitudes.

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Introduction. Lattice modulation spectroscopy has advanced to a standard technique in the physics of ultracold atoms in an optical lattice [1–17]. Of particular interest is the possibility to measure the Mott gap and the creation and analysis of long-lived doublons in a fermionic Hubbard model with strong repulsion [2–16].

On the theoretical side, DMRG work has been done in 1D systems [2]. Based on Fermi's golden rule, it has also been possible to utilize an equilibrium theory to estimate correct doublon production rates [15, 16]. Nonequilibrium dynamical mean-field theory calculations have analyzed models with features analogous to lattice modulation spectroscopy by including time-dependent hopping and time-dependent interactions and showing how they affect the double occupancy [8]. In the linear-response limit, quantum Monte-Carlo calculations have studied phase correlations between the double occupancy and the lattice modulation [12].

In a Mott insulator, doubly occupied sites can be interpreted as occupied by "doublon" quasi-particles. Their long life time is due to a separation of energy scales which requires a rather rare high-order (in the hopping) many-body process for a decay to occur [11]. In this Letter, we provide a computational study of the creation of doublons due to lattice depth modulation. We derive time-dependent tight-binding parameters for a modulated lattice and then apply a recently developed computational strong-coupling method [18]. We then validate our approach by making contact with experimental data by Greif et al. [13] who have provided a detailed measurement of the time evolution of the creation process of doublons in a ^{40}K system. Exploring the parameter space further, we find that for sufficiently high modulation am-

plitudes, processes involving the nonlinear combination of two coherent quanta of the many-body system enhanced doublon production rates at a frequency which equals precisely half the value of the Hubbard repulsion. Higher order nonlinear effects are difficult to produce due to the way the amplitude modulation of the potential translates into the time dependence of the microscopic parameters of the single-band Hubbard model.

Method. The time dependence of the lattice depth modulation is set as follows:

$$V(t) = V_0 + \chi_{[0, t_{\text{mod}}]}(t) \cdot \Delta V \sin \omega t, \quad (1)$$

where V_0 , ΔV , and ω are the average value, modulation amplitude, and modulation frequency of the optical lattice potential depth, respectively, and

$$\chi_{[0, t_{\text{mod}}]}(t) = \begin{cases} 1, & \text{if } t \in [0, t_{\text{mod}}], \\ 0, & \text{otherwise} \end{cases} \quad (2)$$

is the characteristic function of the time interval over which the lattice depth is modulated. The modulation time period length $t_{\text{mod}} = n_{\text{mod}} \cdot 2\pi/\omega$ is a function of the number of modulation cycles n_{mod} chosen for the driving of the system.

The Hamiltonian for a single atom in a d -dimensional optical lattice is given by [19]

$$H_{\text{single}}(t) = -\frac{\hbar^2}{2m} \vec{\nabla}^2 + V(t) \cdot \sum_{i=1}^d \sin^2(kx_i), \quad (3)$$

where $k = 2\pi/\lambda$ is the lattice vector, with the laser wavelength λ . A natural energy unit to use is the recoil-energy $E_R = \hbar^2 k^2 / 2m$. Using the respective time-dependent

maximally localized Wannier functions [20], we map the Hamiltonian to a single-band lattice model. Note that for certain frequencies and amplitudes, transitions to higher bands will eventually become important. For inter-band transitions in particular, corrections from terms coming from the time-derivative of the Wannier functions have to be taken into account [17]. But this is not needed for the case we evaluate here, as we always keep the system in the single-band limit.

The many-body physics of fermionic atoms with spin $1/2$ is then described by the single-band fermionic Hubbard model [21]

$$H(t) = -J(t) \sum_{\langle i,j \rangle, \sigma} (c_{i\sigma}^\dagger c_{j\sigma} + \text{h.c.}) + U(t) \sum_i n_{i\uparrow} n_{i\downarrow} + \sum_{i\sigma} \epsilon(t) n_{i\sigma}, \quad (4)$$

where $J(t) = -\langle w_i(t) | H_{\text{single}}(t) | w_j(t) \rangle$ is the hopping between Wannier states $w_i(t)$ and $w_j(t)$ at neighboring sites i and j , $\epsilon(t) = \tilde{\epsilon}(t) - U(t)/2 - \mu$ is the on-site energy with $\tilde{\epsilon}(t) = \langle w_i(t) | H_{\text{single}}(t) | w_i(t) \rangle$, and $U(t) = g \int d^d r |w_i(\vec{r}; t)|^4$ is the time-dependent repulsion of atoms, while $g = 4\pi\hbar^2 a/m$ is determined by the s -wave scattering length a [6]. The bracket $\langle \cdot, \cdot \rangle$ denotes nearest-neighbor pairs on the lattice. Note that the assumed time-dependence of the single-particle energies is unimportant, since in this work, we perform calculations on a homogeneous, translationally invariant lattice with a fixed density of fermions, and in this case equal-time expectation values are independent of the time evolution of $\epsilon(t)$. The comparison to experiment is made within the Mott insulating region, where we set the chemical potential to $\mu = \tilde{\epsilon}(0)$, for half filling.

In order to compute non-equilibrium observables as a function of time we use a strong-coupling approach which self-consistently expands the self-energy to second order in the hopping [18]. The formalism enables us to numerically evaluate the on-site contour-ordered Green's function

$$G_\sigma(t, t') = -i \langle T_{\mathcal{C}} c_\sigma(t) c_\sigma^\dagger(t') \rangle, \quad (5)$$

where t and t' are times located on the Kadanoff-Baym-Keldysh contour \mathcal{C} [18]. In order to evaluate the site's double occupancy $D(t) = \langle n_\uparrow n_\downarrow \rangle(t)$, we use the following relation for its equal-time derivative, where $t <_{\mathcal{C}} t'$:

$$\left. \frac{\partial G_\sigma(t, t')}{\partial t} \right|_{t'=t+} = U(t) D(t) + \epsilon(t) \langle n_\sigma \rangle(t) + e_\sigma^{\text{kin}}(t). \quad (6)$$

In this expression, the contribution of the spin state σ to the kinetic energy per atom can be evaluated in momentum space via

$$e_\sigma^{\text{kin}}(t) = -2 \frac{J(t)}{N_k} \sum_k^{N_k} \langle n_{k\sigma} \rangle(t) \sum_m^d \cos k_m. \quad (7)$$

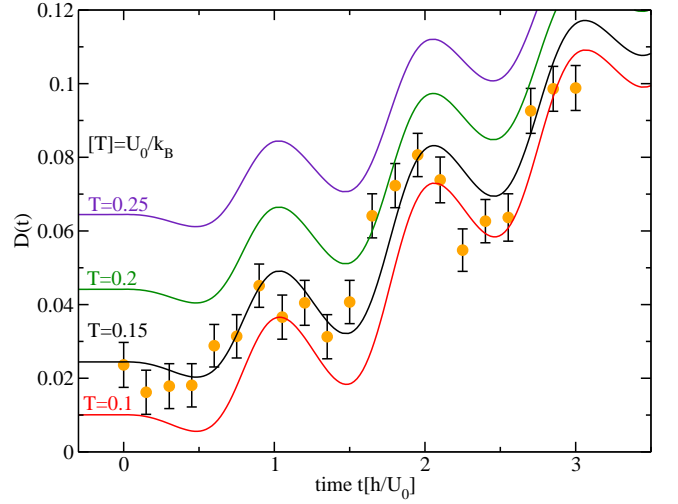


FIG. 1: (color online) Comparison of the double occupancy as a function of time to the experimental data from Ref. [13]. Given parameters of the experiment are $U_0/6J_0 = 4.1$, $V_0 = 7E_R$, $\Delta V/V_0 = 20\%$, $\hbar\omega = U_0$. Since the initial temperature T is the only unknown parameter, we plotted results of our theory for several values of T . The temperature decreases from top to bottom as indicated by the labels. Error bars denote statistical errors from multiple measurements.

Comparison to Experiment. In order to validate the strong-coupling approach, Fig. 1 provides a comparison to recent experimental data on the process of doublon creation due to lattice modulation in a 3D Hubbard model with ^{40}K [13]. The figure shows the double occupancy $D(t)$ as a function of time as the lattice depth $V(t)$ is modulated. The initial value of $V(t)$ is $V_0 = 7E_R$, and it is modulated by $\Delta V/V_0 = 20\%$ at a frequency $\hbar\omega = U_0$. The two-body scattering length is tuned through a Feshbach resonance such that $U_0/6J_0 = 4.1$. For an initial temperature $k_B T = 0.15U_0$, the double occupancy matches the initial experimental value at $t = 0$, and the theory yields a good description of the subsequent time-dependence. Since the lattice modulation frequency equals the Hubbard repulsion $U_0 = U(0)$, particles are resonantly excited from the lower to upper Hubbard band during the lattice modulation. There also exists a process of de-excitation which is represented by the decreasing sections of the observed curve. For the system studied in Fig. 1, the excitation process dominates.

Fig. 2 shows the same analysis for different modulation frequencies. Clearly, the doublon production rates at these frequencies are much lower than for the resonant case $\hbar\omega = U_0$. As a consequence, the signal-to-noise ratio in the experiment is also higher. We again find that $k_B T \approx 0.15U_0$ is a reasonable estimate for the initial temperature. In some cases, however, the data seem to be more compatible with $k_B T = 0.1U_0$ or $k_B T = 0.07U_0$. Although the experimental preparation procedure of the initial thermal state was identical for each data point,

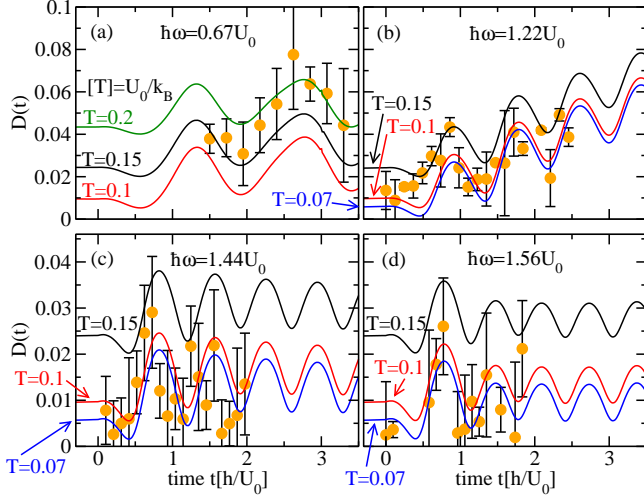


FIG. 2: (color online) Same as Fig. 1 for off-resonant modulation frequencies. Data provided by the Esslinger group. The temperature decreases from top to bottom in each panel. Error bars here are just the spreads from three successive measurements.

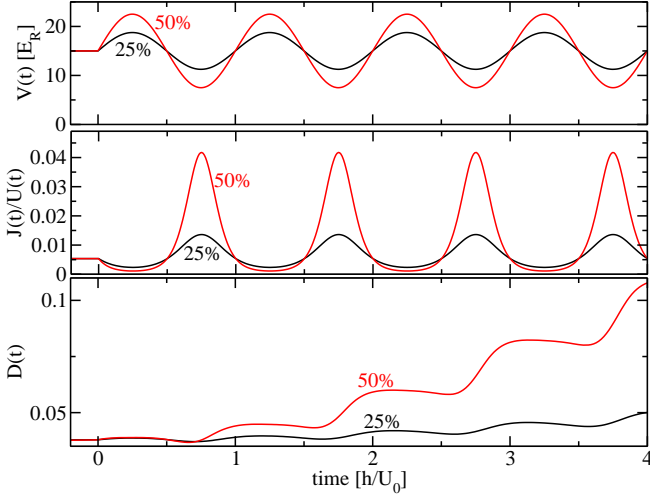


FIG. 3: (color online) Effect of the amplitude on lattice parameters and doublon production. Model parameters are $V_0 = 15E_R$, $U_0/6J_0 = 31.10$, $k_B T = 0.2U_0$, $\hbar\omega = U_0$. The modulation amplitude is either 25% or 50%, as specified for each graph.

systematic drifts in the temperature are quite possible and hard to predict or measure.

Amplitude Effects. Let us now discuss the effect of the amplitude ΔV in more detail. In order to prevent inter-band transitions from becoming important, we study this problem for a deeper lattice, such that $\min_{t \in [0, t_{\text{mod}}]} V(t) \geq 7E_R$. We have plotted the lattice depth $V(t)$, the normalized hopping $J(t)/U(t)$, and the double occupancy $D(t)$ as a function of time for a deeper lattice at two strong values of the lattice modulation am-

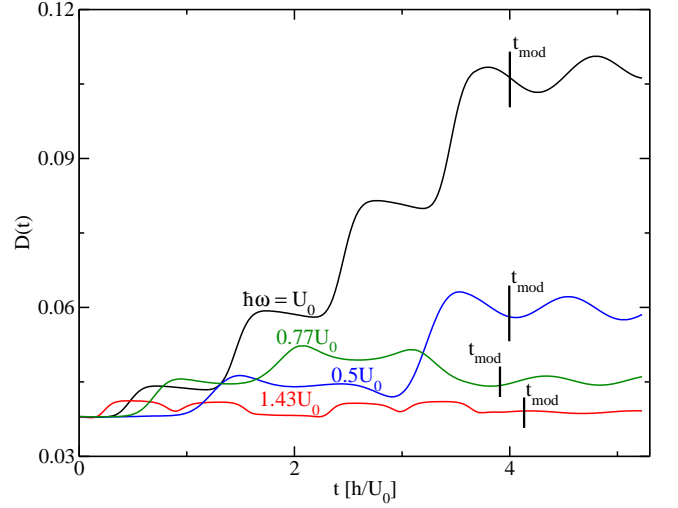


FIG. 4: (color online) Double occupancy as a function of time for lattice modulations at different frequencies. The other model parameters are same as in Fig. 3 for $\Delta V/V_0 = 50\%$. Vertical bars denote the respective end time of the modulation. The modulation frequency associated with a given graph is shown with a label on the curve.

plitude in Fig. 3. As the amplitude is raised, the non-linear relationship between the hopping and the lattice depth results in a periodically kicked rather than a periodically driven system (due to $J(t)/U(t)$ becoming very small for deep lattices). Overall, the double occupancy is increased stepwise within each modulation cycle. Again, we observe both, excitation and de-excitation processes in the double occupancy data. Within the first modulation cycle, we observe the former and the latter to coincide with the decreasing and increasing regimes of $J(t)/U(t)$, respectively. Within modulation cycles further out in time, the stepwise increase in double occupancy starts already when $J(t)/U(t)$ is still rising. This is due to the fact that between the spikes in $J(t)/U(t)$, the hopping is effectively zero, so that the system oscillates internally with frequency $U(t) \approx U_0 = \hbar\omega$ within the four-level Hilbert space associated with a single lattice site [22]. As a consequence, there is a constructive interference of the immediate effect of lattice modulation and the internal oscillation, once the latter is fully engaged.

Amplitude versus Frequency. We would like to address the interplay of the internal oscillation at zero hopping and the lattice modulation in more detail next. Let us discuss data at different frequencies by first assuming a strong lattice modulation amplitude $\Delta V/V_0 = 50\%$. The modulation time interval in Eq. (1) for a given frequency is chosen to be $t_{\text{mod}} := \left\lfloor \frac{4.93\hbar/U_0}{2\pi/\omega} \right\rfloor \cdot \frac{2\pi}{\omega}$. Results are shown in Fig. 4. The strongest increase in double occupancy is observed in the resonant case $\hbar\omega = U_0$. As the end time t_{mod} of the modulation is exceeded, the system contin-

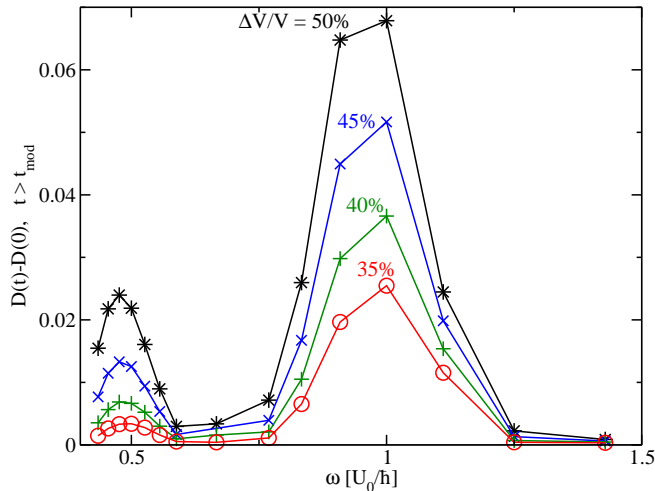


FIG. 5: (color online) Doubloon production as a function of frequency for different amplitudes. Amplitude increases from bottom to top. The other parameters are the same as in Figs. 3 and 4.

ues oscillating with the internal frequency U_0 . A significant increase in double occupancy is also observed in the case $\hbar\omega = U_0/2$. Moreover, in this case also the system continues oscillating with U_0 after modulation, but with a smaller amplitude. For the off-resonant frequencies $\hbar\omega = 1.43U_0$ and $\hbar\omega = 0.77U_0$, some intermediate excitations to the upper Hubbard band show up but (partially) annihilate after a while. The amplitude of internal oscillations after the modulation is turned off appears to grow with the total increase in double occupancy reached when the modulation ends.

To provide a further overview of the frequency and amplitude dependence, Fig. 5 shows the final value of the double occupancy, i.e. the value averaged over one oscillation cycle $[t_{\text{mod}}, t_{\text{mod}} + h/U_0]$, vs. the frequency for different amplitudes as a lattice modulation “spectroscopy”. Evidently, a second-order resonance is observed at frequency $\hbar\omega = 0.5U_0$, which is suppressed as the amplitude is lowered. It can be interpreted as a coherent excitation involving the nonlinear combination of two smaller quanta of energy $U_0/2$, i.e. an analogue to two-photon excitations in quantum optics. As a consequence of the combination of these two quanta, the width of the second-order peak is approximately half of the width of the first-order peak. Also note that both peaks are shifted towards slightly smaller frequencies due to the finite width of the Hubbard bands. Lower-energy excitations are possible by exciting from the upper edge of the lower Hubbard band to the lower edge of upper Hubbard band.

Whereas the $U_0/2$ peak is clearly visible for very large modulation amplitudes, the rather isolated time evolution of lattice sites between the kicks in $J(t)/U(t)$ presumably suppresses its amplitude. This is inherent to

modulation spectroscopy due to the non-linearity of the map $V \rightarrow J/U$. Its presence could be possibly enhanced by designing $V(t)$ in such a way that a harmonic shape is obtained for $J(t)$.

Summary. We have studied finite-amplitude lattice depth modulation spectroscopy of ultracold fermionic atoms in the Mott-insulating phase using a recently developed strong-coupling method [18]. In order to validate the theory, we have compared to experimental data and found excellent agreement. Only the temperature of the initial thermal state was unknown in the experiment and had to be determined *a posteriori*. We furthermore analyzed higher amplitudes of the modulation strength for deeper lattices. A large value of the amplitude results in a pulsed, rather than a driven system, in terms of the time-dependence of the hopping. This causes step-shaped changes in the double occupancy, accompanied only by oscillations with frequency $U(t) \approx U_0$ of local degrees of freedom on a single lattice site. At a certain threshold in amplitude, a second “nonlinear” peak in the doublon production rate appears at $\omega = U_0/2$.

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* Electronic address: andreas@physics.georgetown.edu

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- [22] In this nonlinear regime, the average of $U(t)$ becomes slightly smaller than U_0 . The variation in $U(t)$ is much smaller than that of $J(t)$, when the amplitude of the modulation is 50%. For example, when in units of J_0 , $J(t)$ varies by more than a 100%, since it is nearly completely suppressed when $V(t)$ is large and increased by up to 400% when $V(t)$ is small, while $U(t)$ varies by only 30% of U_0 .